

# Math 307 Lecture 5

## Modeling Equations like a Pro!

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# Today!

Last time:

- Review of first order linear and separable ODEs

This time:

- First homework due Friday!!!
- Modeling first-order equations

Next time:

- Differences between linear and nonlinear equations

# Outline

- 1 Linear Equation for Mixing
- 2 Radiative Heat Transfer
- 3 Holes in a Bucket
  - Torricelli's Law
  - A story problem
- 4 Pebble falling in Syrup
  - Stoke's Law
  - A story problem

## Our First Example Model: Mixing fluids

Figure : Rate of pollution of a pond can be modeled by a linear ordinary differential equation



- Polluted water flows into a pond
- Volume of pond (constant):  $V = 10^7$  gal
- Amount of pollutant in pond:  $P$  (metric tons)
- Toxic sludge flows in at  $5 \times 10^6$  gal/yr
- Toxic sludge contains  $2 + \sin(2t)$  grams of pollutant per gallon
- Lake unpolluted at  $t = 0$

## Our First Example Model: Mixing fluids

### Question

How does the amount of pollutant change over time?

$$\frac{dP}{dt} = \text{rate in} - \text{rate out}$$

$$\text{rate in} = \overbrace{(5 \times 10^6)}^{\text{gal. toxic sludge/yr}} \cdot \underbrace{(2 + \sin(2t))}_{\text{grams pollutant/gal}} \cdot \overbrace{(10^{-6})}^{\text{grams/ton}}$$

$$\text{rate out} = \overbrace{(5 \times 10^6)}^{\text{gal. mixed pond water/yr}} \cdot \underbrace{P(t)/V}_{\text{metric tons pollutant/gal}}$$

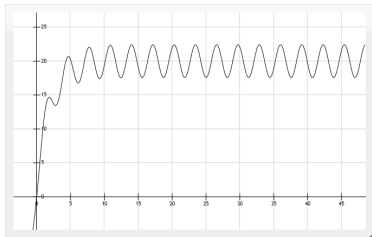
## Our First Example Model: Mixing fluids

- IVP  $P(0) = 0$  and

$$\frac{dP}{dt} = 10 + 5 \sin(2t) - \frac{1}{2}P$$

- Int. factor is  $\mu(t) = e^{t/2}$
- Solution given below
- Limit behavior:  
oscillation about  $P=20$

Figure : Pond pollution modeled by a linear ODE



$$P(t) = 20 - \frac{40}{17} \cos(2t) + \frac{10}{17} \sin(2t) - \frac{300}{17} e^{-t/2}$$

## Our Next Example Model: Radiative Heat Transfer

Figure : Radiative heat



- $U$  abs. temp. of body
- $T$  abs. temp. of space
- $\alpha$  emissivity constant

- Stefan-Boltzmann law:

$$\frac{dU}{dt} = -\alpha(U^4 - T^4)$$

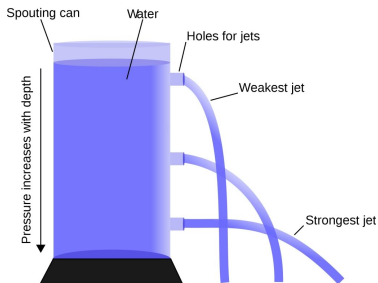
- For  $U \gg T$ , we can approximate

$$\frac{dU}{dt} = -\alpha U^4$$

- Separable! Solution is  
$$U = \frac{1}{(3\alpha t + C)^{1/3}}$$

# Torricelli's Law

Figure : Water under more pressure shoots faster/farther



- Torricelli's law says
- $v = \sqrt{2gh}$
- Outflow velocity:  $v$
- Water level above opening:  $h$
- How fast does water leave the tank?



## An Old West Example

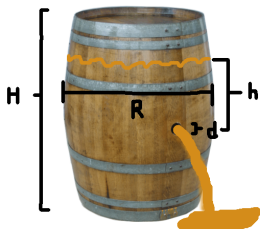
**Figure :** “You’ve got to ask yourself one question: ‘Do I feel lucky?’ Well, do ya punk?” – Clint Eastwood on Differential Equations



- Clint Eastwood shoots a cylindrical barrel of whiskey dead center
- bullet does not leave barrel
- barrel has height 3 feet and radius 1 foot
- he uses a .44 Magnum, the most powerful handgun in the world
- how long until the barrel is half-empty?

# An Old West Example

Figure : "Glub glub glub..."



- $H = 3$  ft,  $R = 1$  ft
- $d = 10.9$  millimeters  
(from size of bullet)
- $V$  volume of barrell
- area of hole:  $A = \pi d/2 = 0.297 \cdot 10^{-4} \pi$  ft
- $dV/dt = -A \cdot v_{out}$

## An Old West Example

- Toricelli says  $v_{\text{out}} = \sqrt{2gh}$ .
- Therefore

$$\frac{dV}{dt} = -A\sqrt{2gh}$$

- Also since  $V = \pi R^2 h$ , we have

$$\frac{dV}{dt} = \pi R^2 \frac{dh}{dt}$$

- Therefore

$$-A\sqrt{2gh} = \pi R^2 \frac{dh}{dt}$$

## An Old West Example

- Separable equation!

$$\frac{1}{\sqrt{h}} \frac{dh}{dt} = -\frac{A\sqrt{2g}}{\pi R^2}$$

$$2\sqrt{h} = -\frac{A\sqrt{2g}}{\pi R^2}t + C$$

$$h = \left( -\frac{A\sqrt{g}}{\sqrt{2}\pi R^2}t + C \right)^2$$

- At  $t = 0$ ,  $h = H/2$ , so  $C = \sqrt{H/2}$  and

$$h = \frac{1}{2} \left( \sqrt{H} - \frac{A\sqrt{g}}{\pi R^2}t \right)^2$$

## An Old West Example

- Set  $h = 0$  and solve for  $t$ :

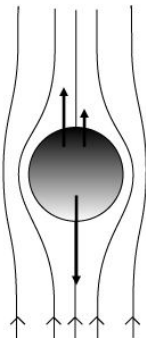
$$t = \frac{\sqrt{H}\pi R^2}{A\sqrt{g}}$$

- Putting in the values of  $A, H, g$  ( $g = 32 \text{ ft/s}^2$ )

$$t = 8.42 \cdot 10^3 \text{ s} \approx 14 \text{ min.}$$

# Stokes's Law

Figure : Falling pebble feels three forces



- Stokes law governs the drag felt by an object falling through a viscous fluid
- Spherical pebble of radius  $r$ , mass  $m$ , and velocity  $v$
- Ball feels three forces: buoyant force, gravitational force, viscous drag

# Stokes's Law

**Figure :** Stokes was known for being super stoked about fluid mechanics



- buoyant force:  $B$  equals weight of fluid displaced
- viscous drag:  
 $R = -6\pi\mu rv$  by Stokes law
- here  $\mu$  quantifies how viscous the fluid is
- $\mu$  is bigger for molasses than for water
- gravitational force:  $-mg$
- How fast does a pebble sink?

# 3000 Leagues under the Sea

Figure : The Beatles did math in a yellow submarine...maybe



- We drop a spherical submarine into the ocean
- assume water density is approximately constant with respect to depth
- buoyant force is then
$$B = \frac{4}{3}\pi r^3 \rho g$$
- by Newtons law:
$$F = m \frac{dv}{dt}$$
- total force

$$F = -mg + B + R$$



## 3000 Leagues under the Sea

- velocity therefore satisfies the differential equation

$$m \frac{dv}{dt} = \overbrace{-mg}^{\text{grav. force}} + \underbrace{\frac{4}{3}\pi r^3 \rho g}_{\text{buoyant force}} - \overbrace{6\pi\mu r v}^{\text{viscous drag}}$$

- Linear equation! (alt. it's separable)
- Find an integrating factor and solve. Should get:

$$v = C \exp\left(\frac{-6\pi\mu r}{m} t\right) + \frac{-mg + \frac{4}{3}\pi r^3 \rho g}{6\pi\mu r}$$

# 3000 Leagues under the Sea

Figure : Sea monsters love differential equations



- as we fall, we gather more speed!
- what is the terminal velocity of the submarine (maximum speed it can fall)?
- assume not attacked by a sea monster
- take limit as  $t \rightarrow \infty$

$$v_{\text{term}} = \frac{-mg + \frac{4}{3}\pi r^3 \rho g}{6\pi \mu r}$$

# Summary!

What we did today:

- We looked at some real-world situations that can be modeled by differential equations

Plan for next time:

- Differences between linear and nonlinear equations