

MATH 309: Homework #4

Due on: November 20, 2015

Problem 1 *Even and Odd Functions*

Prove that any function $f(x)$ may be expressed as a sum of two functions $f(x) = g(x) + h(x)$ with $g(x)$ even and $h(x)$ odd. [Hint: consider $f(x) + f(-x)$].

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Solution 1. In order to prove the statement we want, we need to show that for any function $f(x)$, there exists an even function $g(x)$ and an odd function $h(x)$ with $f(x) = g(x) + h(x)$. In particular, we need to come up with equations for $g(x)$ and $h(x)$ in terms of $f(x)$. How can we do this? One way is to assume that $g(x)$ and $h(x)$ are known to exist, and then fiddle around with $f(x)$ to figure out the equations. In particular if $g(x)$ is even and $h(x)$ is odd and $f(x) = g(x) + h(x)$ then

$$f(-x) = g(-x) + h(-x) = g(x) - h(x).$$

It follows that

$$f(x) + f(-x) = g(x) + h(x) + (g(x) - h(x)) = 2g(x),$$

and therefore we should take $g(x) = (f(x) + f(-x))/2$. Similarly, we have that

$$f(x) - f(-x) = g(x) + h(x) - (g(x) - h(x)) = 2h(x),$$

and therefore we should take $h(x) = (f(x) - f(-x))/2$. Great!

What we did above is just a bunch of scratch work. Here's our actual proof: Suppose that $f(x)$ is a function. Define $g(x) = (f(x) + f(-x))/2$ and $h(x) = (f(x) - f(-x))/2$. Then since

$$g(-x) = (f(-x) + f(-(-x)))/2 = (f(-x) + f(x))/2 = (f(x) + f(-x))/2 = g(x)$$

we have that $g(x)$ is even. Similarly

$$h(-x) = (f(-x) - f(-(-x)))/2 = (f(-x) - f(x))/2 = -(f(x) - f(-x))/2 = -h(x)$$

and therefore $h(x)$ is odd. Furthermore

$$g(x) + h(x) = (f(x) + f(-x))/2 + (f(x) - f(-x))/2 = f(x).$$

Therefore $f(x) = g(x) + h(x)$ is a sum of an even function and an odd function. This completes our proof.

Problem 2 *Even and Odd Functions*

Prove that the derivative of an even function is odd and that the derivative of an odd function is even.

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Solution 2. There are many great ways to prove this fact. We will use one of the simplest methods: the chain rule. Let $g(x) = f(-x)$. Then by the chain rule

$$g'(x) = -f'(-x).$$

Now let's suppose $f(x)$ is an even function. Then in this case $g(x) = f(x)$, making $g'(x) = f'(x)$, so that the above expression reads $f'(x) = -f'(-x)$. Since x was arbitrary, this shows that $f'(x)$ is odd when $f(x)$ is even. Alternatively, let's suppose that $f(x)$ is an odd function. Then $g(x) = -f(x)$, making $g'(x) = -f'(x)$, so that the expression we derived from the chain rule reads $-f'(x) = -f'(-x)$, and hence $f'(x) = f'(-x)$. Since x was arbitrary, this shows that $f'(x)$ is even when $f(x)$ is odd. This completes our proof.

Problem 3 *Sine Series*

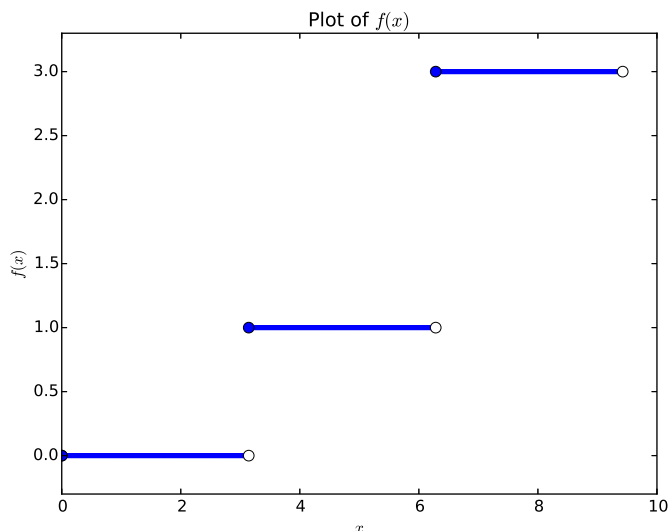
Consider the function

$$f(x) = \begin{cases} 0, & 0 < x < \pi \\ 1, & \pi < x < 2\pi \\ 2, & 2\pi < x < 3\pi \end{cases}$$

- (a) Sketch a graph of $f(x)$
- (b) By reflecting f appropriately, express f as a sine series.
- (c) Plot three different partial sums of the sine series, clearly indicating the partial sums being plotted.
- (d) Sketch a graph of the function to which the sine series converges for three periods.

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Solution 3.



(a)

(b) To express $f(x)$ as a sine series, we create a new function $g(x)$ which is odd and periodic by reflecting $f(x)$ oddly across the y -axis, and then defining $g(x+6\pi) = g(x)$ for all x . Since $g(x)$ is periodic, it has a Fourier series, and since $g(x)$ is odd, all of the cosine terms will be gone, leaving just the sine terms. We can calculate the associated coefficients by using the Euler-Fourier formula:

$$b_n = \frac{1}{3\pi} \int_{-3\pi}^{3\pi} g(x) \sin(n\pi x / (3\pi)) dx.$$

Now since $g(x)$ is odd, the integrand is even, so we can simply integrate from 0 to 3π and multiply by 2 to get the value of b_n . Moreover, from 0 to 3π the function $g(x)$ agrees with $f(x)$, and therefore

$$b_n = \frac{2}{3\pi} \int_0^{3\pi} f(x) \sin(nx/3) dx.$$

Now in order to do this integral, we need to break it up into the three separate intervals where $f(x)$ is individually defined:

$$b_n = \frac{2}{3\pi} \left(\int_0^\pi 0 \sin(nx/3) + \int_\pi^{2\pi} 1 \sin(nx/3) dx + \int_{2\pi}^{3\pi} 3 \sin(nx/3) dx \right).$$

The integrals themselves are pretty easy. Evaluating them, we obtain:

$$b_n = \frac{2}{3\pi} \left(0 + \frac{-3}{n} (\cos(2n\pi/3) - \cos(n\pi/3)) dx + \frac{-9}{n} (\cos(3n\pi/3) - \cos(2n\pi/3)) \right).$$

Now we want to use the fact that

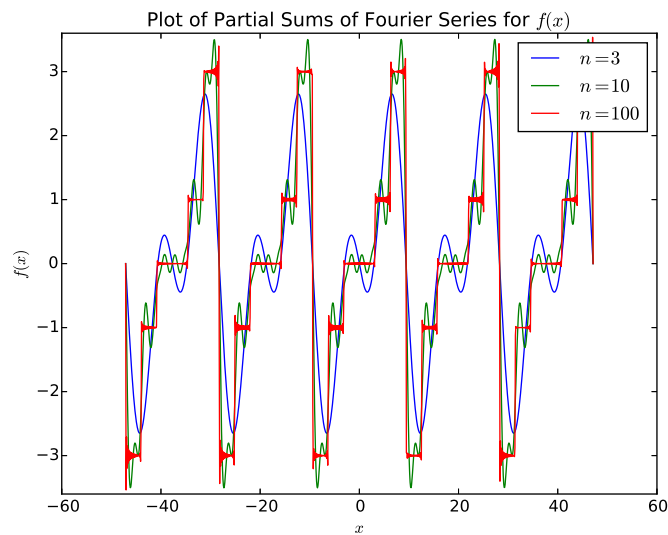
$$\cos(m\pi/3) = \begin{cases} 1/2, & m = \pm 1 + 6k \\ -1/2, & m = \pm 2 + 6k \\ 1, & m = 0 + 6k \\ -1, & m = 3 + 6k \end{cases}$$

Using this, the expression for b_n reduces to

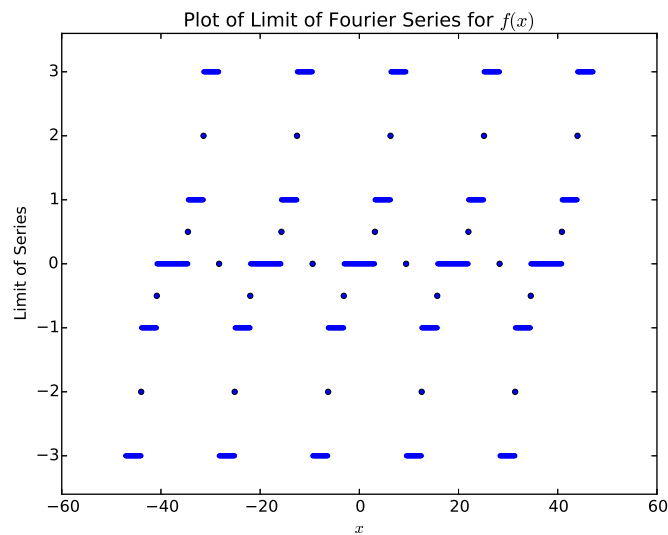
$$b_n = \begin{cases} 5/(n\pi), & n = \pm 1 + 6k \\ -9/(n\pi), & n = \pm 2 + 6k \\ 0, & n = 0 + 6k \\ 8/(n\pi), & n = 3 + 6k \end{cases}$$

Using these values of b_n , we have

$$f(x) = \sum_{n=1}^{\infty} b_n \sin(nx/3).$$



(c)



(d)

Problem 4 *Cosine Series*

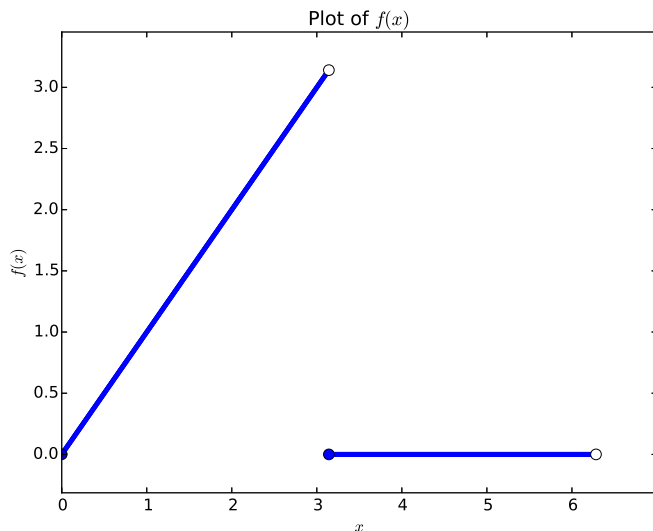
Consider the function

$$f(x) = \begin{cases} x, & 0 < x < \pi \\ 0, & \pi < x < 2\pi \end{cases}$$

- (a) Sketch a graph of $f(x)$
- (b) By reflecting f appropriately, express f as a cosine series.
- (c) Plot three different partial sums of the cosine series, clearly indicating the partial sums being plotted.
- (d) Sketch a graph of the function to which the cosine series converges for three periods.

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Solution 4.



(a)

- (b) To express $f(x)$ as a cosine series, we create a new function $g(x)$ which is even and periodic by reflecting $f(x)$ evenly across the y -axis, and then defining $g(x+4\pi) = g(x)$ for all x . Since $g(x)$ is periodic, it has a Fourier series, and since $g(x)$ is even, all of the sine terms will be gone, leaving just the cosine terms. We can calculate the associated coefficients by using the Euler-Fourier formula:

$$a_n = \frac{1}{2\pi} \int_{-2\pi}^{2\pi} g(x) \cos(n\pi x / (2\pi)) dx.$$

Now since $g(x)$ is odd, the integrand is even, so we can simply integrate from 0 to 2π and multiply by 2 to get the value of a_n . Moreover, from 0 to 2π the function

$g(x)$ agrees with $f(x)$, and therefore

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos(nx/2) dx.$$

Now in order to do this integral, we need to break it up into the two separate intervals where $f(x)$ is individually defined:

$$a_n = \frac{1}{\pi} \left(\int_0^\pi x \cos(nx/2) + \int_\pi^{2\pi} 0 \cos(nx/2) dx \right).$$

To evaluate this integral, we use integration by parts, obtaining:

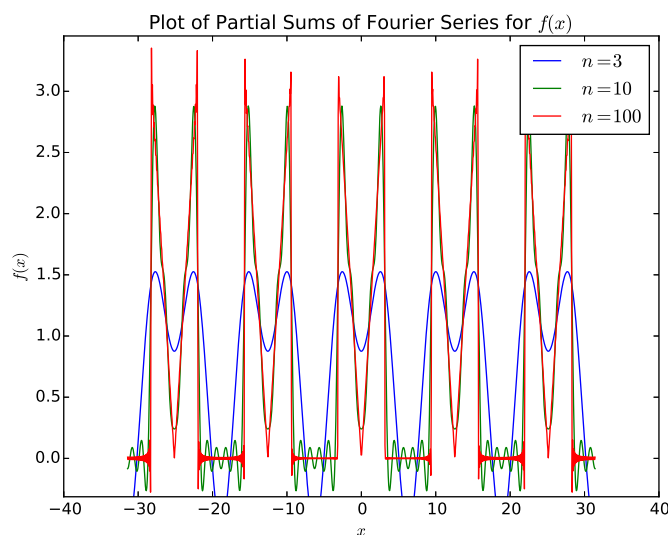
$$a_n = \frac{-2}{n} \cos(n\pi/2) = \begin{cases} ((-1)^{n/2} - 1)4/(n^2\pi) & n \text{ even} \\ (-1)^{(n+1)/2}2/n - 4/(n^2\pi) & n \text{ odd} \end{cases}$$

This expression does not work however for $n = 0$ since in the calculation we divided by n . We must do this separately:

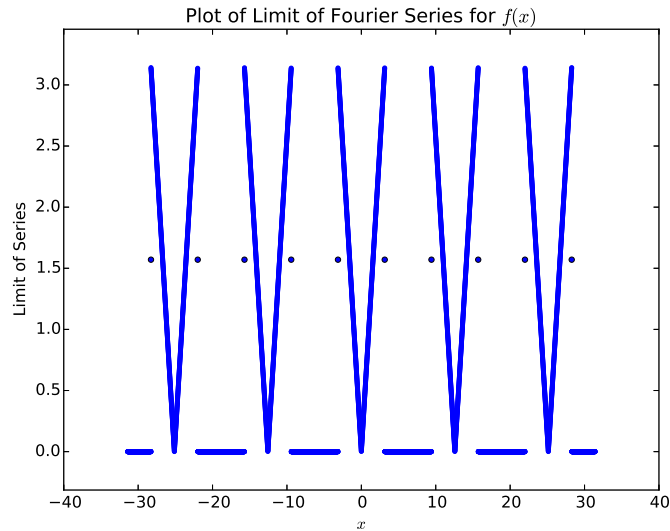
$$a_0 = \frac{1}{\pi} \int_0^{2\pi} f(x) dx = \frac{1}{\pi} \int_0^\pi x dx = \frac{1}{2}\pi.$$

Using these values of a_n , we have

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(n\pi x/3).$$



(c)



(d)

Problem 5 *Heat Equation 1*

Find the solution of the heat conduction problem

$$\begin{aligned}
 100u_{xx} &= u_t, & 0 < x < 1, & t > 0 \\
 u(0, t) &= u(1, t) = 0, & t > 0 \\
 u(x, 0) &= \sin(2\pi x) - \sin(5\pi x)
 \end{aligned}$$

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Solution 5. We identify $\alpha^2 = 100$ and $L = 1$. Then we need to expand $u(x, 0)$ as

$$u(x, 0) = \sum_{n=1}^{\infty} b_n \sin(n\pi x).$$

However, if we look at the form of $u(x, 0)$, this is immediately accomplished by taking $b_2 = 1, b_5 = -1$ and $b_n = 0$ otherwise. Therefore

$$u(x, t) = u_2(x, t) - u_5(x, t) = e^{-400\pi^2 t} \sin(2\pi x) - e^{-2500\pi^2 t} \sin(5\pi x).$$

Problem 6 *Heat Equation 2*

Find the solution of the heat conduction problem

$$\begin{aligned}
 u_{xx} &= 4u_t, & 0 < x < 2, & t > 0 \\
 u(0, t) &= u(2, t) = 0, & t > 0 \\
 u(x, 0) &= 2 \sin(\pi x/2) - \sin(\pi x) + 4 \sin(2\pi x)
 \end{aligned}$$

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Solution 6. We identify $\alpha^2 = 4$ and $L = 2$. Then we need to expand $u(x, 0)$ as

$$u(x, 0) = \sum_{n=1}^{\infty} b_n \sin(n\pi x/2).$$

However, if we look at the form of $u(x, 0)$, this is immediately accomplished by taking $b_1 = 2, b_2 = -1, b_4 = 4$ and $b_n = 0$ otherwise. Therefore

$$u(x, t) = u_1(x, t) - u_2(x, t) + 4u_4(x, t) = e^{-\pi^2 t} \sin(\pi x/2) - e^{-4\pi^2 t} \sin(\pi x) + 4e^{-16\pi^2 t} \sin(2\pi x).$$

Problem 7 *Schrödinger Equation*

In quantum mechanics, the position of a point particle in space is not certain – it's described by a probability distribution. The probability distribution of the position of the particle is $|\psi(x, t)|^2$, where $\psi(x, t)$ is the **wave function** of the particle. (Note: the wave function $\psi(x, t)$ can be complex-valued!!). The one-dimensional, time-dependent Schrödinger equation, describing the wave function $\psi(x, t)$ of a particle of mass m interacting with a potential $v(x)$ is given by

$$i\hbar\psi_t(x, t) = -\frac{\hbar^2}{2m}\psi_{xx}(x, t) + v(x)\psi(x, t)$$

where \hbar is some universal constant. The potential $v(x)$ can be imagined as a function describing the particles interaction with whatever “stuff” is in the space surrounding the particle, eg. walls, external forces, etc.

- (a) Use separation of variables to replace this partial differential equation with a pair of two ordinary differential equations
- (b) If $v(x)$ is a potential corresponding to an “infinite square well”:

$$v(x) = \begin{cases} 0, & -1 < x < 1 \\ \infty, & |x| \geq 1 \end{cases}$$

Then $\psi(x, t)$ must be zero whenever $|x| \geq 1$ and therefore $\psi(x, t)$ is the wave function of a particle trapped in a one-dimensional box! In other words, this potential describes a particle surrounded by impermeable walls. In this case, Schrödinger's equation reduces to

$$\begin{aligned} i\hbar\psi_t(x, t) &= -\frac{\hbar^2}{2m}\psi_{xx}(x, t), & -1 < x < 1, & t > 0 \\ \psi(-1, t) &= \psi(1, t) = 0, & t > 0 \end{aligned}$$

Suppose that initially the wave function is known to be

$$\psi(x, 0) = \frac{3}{5} \sin(\pi x) + \frac{4}{5} \sin(3\pi x).$$

Determine $\psi(x, t)$ for all $t > 0$.

- (c) Since $|\psi(x, t)|^2$ is the probability *distribution* of the particle's position at time t , the probability that the particle is somewhere in the box between ℓ_1 and ℓ_2 is given by

$$\mathbb{P}(\ell_1 \leq \text{pos} \leq \ell_2) = \int_{\ell_1}^{\ell_2} |\psi(x, t)|^2 dx.$$

Show that the probability $\mathbb{P}(-1 \leq \text{pos} \leq 1)$ that the particle is between -1 and 1 is always 1 (in other words, the particle is always in the box!).

- (d) What is the probability $\mathbb{P}(-1 \leq \text{pos} \leq 0)$ that the particle is in the first half of the box at any given time?

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Solution 7.

- (a) We assume $\psi(x, t) = F(x)G(t)$. Then inserting this into Schrödinger's equation, we obtain

$$i\hbar F(x)G'(t) = -\frac{\hbar^2}{2m}F''(x)G(t) + v(x)F(x)G(t).$$

Now if we divide out by a $G(t)$ and a $F(x)$ we find

$$i\hbar G'(t)/G(t) = -\frac{\hbar^2}{2m}F''(x)/F(x) + v(x).$$

The function on the left hand side is a function of t only. The function on the right hand side is a function of x only. Therefore the only way that the above equality can work is if both sides are equal to some constant E . Therefore

$$i\hbar G'(t)/G(t) = E, \quad -\frac{\hbar^2}{2m}F''(x)/F(x) + v(x) = E.$$

Simplifying, this gives us the system of two ordinary differential equations

$$i\hbar G'(t) = EG(t).$$

$$-\frac{\hbar^2}{2m}F''(x) + v(x)F(x) = EF(x).$$

The latter equation of these two equations is known as the **time-independent Schrödinger equation**.

- (b) This is just like the heat equation, with $\alpha^2 = i\frac{\hbar}{2m}$ and $L = 1$. Thus given the initial condition, the solution that we are looking for is

$$\psi(x, t) = \frac{3}{5}e^{-i\frac{\hbar\pi^2}{2m}t} \sin(\pi x) + \frac{4}{5}e^{-i\frac{9\hbar\pi^2}{2m}t} \sin(3\pi x).$$

(c) Note that

$$\psi(x, t)^* = \frac{3}{5} e^{i\frac{\hbar\pi^2}{2m}t} \sin(\pi x) + \frac{4}{5} e^{i\frac{9\hbar\pi^2}{2m}t} \sin(3\pi x),$$

and therefore

$$|\psi(x, t)|^2 = \psi(x, t)\psi(x, t)^* = \frac{9}{25} \sin^2(\pi x) + \frac{16}{25} \sin^2(3\pi x) + \frac{12}{25} \left(e^{i\frac{8\hbar\pi^2}{2m}t} + e^{i\frac{-8\hbar\pi^2}{2m}t} \right) \sin(\pi x) \sin(3\pi x).$$

If we now integrate over the domain of the box (from -1 to 1), orthogonality tells us the integral of $\sin(\pi x) \sin(3\pi x)$ dies off! Therefore we obtain:

$$\int_{-1}^1 |\psi(x, t)|^2 dx = \frac{9}{25} \int_{-1}^1 \sin^2(\pi x) dx + \frac{16}{25} \int_{-1}^1 \sin^2(3\pi x) dx = \frac{9}{25} + \frac{16}{25} = 1.$$

This shows that the probability that the particle is in the box at any time t is 1 – e.g. it is a certainty.

(d) We can use the work from above to write

$$\int_{-1}^1 |\psi(x, t)|^2 dx = \frac{9}{25} \int_{-1}^0 \sin^2(\pi x) dx + \frac{16}{25} \int_{-1}^0 \sin^2(3\pi x) dx + \frac{12}{25} \left(e^{i\frac{8\hbar\pi^2}{2m}t} + e^{i\frac{-8\hbar\pi^2}{2m}t} \right) \int_{-1}^0 \sin(\pi x) \sin(3\pi x) dx$$

However, since we're not integrating over the full period, we cannot appeal to orthogonality to say that the cross-term dies anymore. However, direct calculation shows that it does indeed die anyway. The sum of the first two integrals is easily calculated to be $1/2$. Therefore the probability that the particle is in the first half of the box at any time t is exactly $1/2$. In other words – at any time the particle is equally likely to be in either side of the box.