

MATH 309: Homework #5

Due on: November 30, 2015

Problem 1 *Insulated Heat Equation Problem*

Consider a uniform rod of length L with an initial temperature given by $u(x, 0) = \sin(\pi x/L)$ with $0 \leq x \leq L$. Assume that both ends of the bar are insulated (this is a homogeneous von Neumann boundary condition for $t > 0$).

- (a) Find the temperature $u(x, t)$. (Note: the initial condition $u(x, 0)$ does not satisfy the boundary conditions, which is fine since we are only asking the boundary conditions to be satisfied for $t > 0$)
- (b) What is the steady state temperature as $t \rightarrow \infty$?
- (c) Let $\alpha^2 = 1$ and $L = 40$. Plot u vs. x for several values of t .

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Solution 1.

- (a) We need to determine the temperature initially in terms of a cosine series. This means reflecting $\sin(\pi x/L)$ *evenly* and then extending periodically. In other words, we're really looking for the cosine series of $|\sin(\pi x/L)|$. Using Euler-Fourier, we obtain

$$a_n = \frac{1}{L} \int_{-L}^L |\sin(\pi x/L)| \cos(n\pi x/L) dx = \frac{2}{L} \int_0^L \sin(\pi x/L) \cos(n\pi x/L) dx.$$

Now in order to complete the last integral on the right, we can adopt several strategies. The most obvious thing is to integrate by parts twice, and then compare sides – however, that is a lot of work. A shorter strategy is to use the addition angle formulas for sine to write:

$$\sin(\pi x/L) \cos(n\pi x/L) = \frac{1}{2}(\sin((1+n)\pi x/L) + \sin((1-n)\pi x/L)).$$

With this in mind, the above integral becomes

$$a_n = \frac{1}{L} \int_0^L (\sin((1+n)\pi x/L) + \sin((1-n)\pi x/L)) dx = \frac{2}{\pi} \left(\frac{1 + (-1)^n}{1 - n^2} \right).$$

However, notice that in our derivation of this formula, we divided by $1 - n$, and therefore the expression we obtained for a_n does not apply when $n = 1$. We must treat this case separately! We calculate using the double angle formula

$$a_1 = \frac{2}{L} \int_0^L \sin(\pi x/L) \cos(\pi x/L) dx = \frac{1}{L} \int_0^L \sin(2\pi x/L) dx = -\frac{1}{2\pi} \cos(2\pi x/L) \Big|_0^L = 0.$$

We conclude that

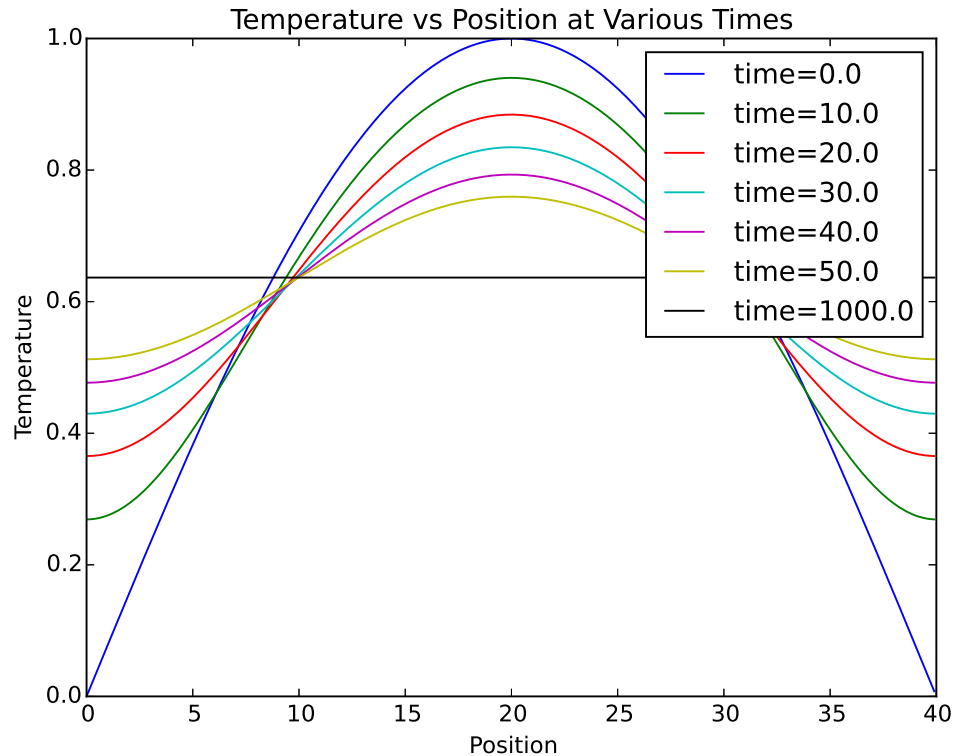
$$u(x, 0) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(n\pi x/L) = \frac{2}{\pi} + \sum_{n=2}^{\infty} \frac{2}{\pi} \left(\frac{1 + (-1)^n}{1 - n^2} \right) \cos(n\pi x/L).$$

This tells us that

$$u(x, t) = \frac{2}{\pi} + \sum_{n=2}^{\infty} \frac{2}{\pi} \left(\frac{1 + (-1)^n}{1 - n^2} \right) e^{-n^2 \pi^2 \alpha^2 t / L^2} \cos(n\pi x/L).$$

(b) As $t \rightarrow \infty$, the exponential terms die off, leaving only $a_0/2$. Therefore the steady state temperature is $2/\pi$.

(c) Plot at several times is included in the figure below.



Problem 2 *Another Insulated Heat Equation Problem*

Consider a bar of length 40 cm whose initial temperature is given by $u(x, 0) = x(60 - x)/30$. Suppose that $\alpha^2 = 1/4$ cm²/s and that both ends of the bar are insulated.

- (a) Find the temperature $u(x, t)$. (Note: the initial condition $u(x, 0)$ does not satisfy the boundary conditions, which is fine since we are only asking the boundary conditions to be satisfied for $t > 0$)
- (b) What is the steady state temperature as $t \rightarrow \infty$?
- (c) Plot u vs. x for several values of t .
- (d) Determine how much time must elapse before the temperature at $x = 40$ comes within 1 degrees C of its steady state value.

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Solution 2.

- (a) Again, we must extend $u(x, 0)$ evenly and periodically in order to pick up its cosine series. Then by the Euler-Fourier equation we have

$$a_n = \frac{2}{40} \int_0^{40} \frac{x(60 - x)}{30} \cos(n\pi x/40) dx.$$

We can obtain the value explicitly by using integration by parts twice to get the a'_n 's. (There are, of course, more clever ways to do things, but this works fine). Doing so, we obtain

$$a_n = \frac{160}{3} \frac{(-1)^{n+1} - 3}{n^2\pi^2},$$

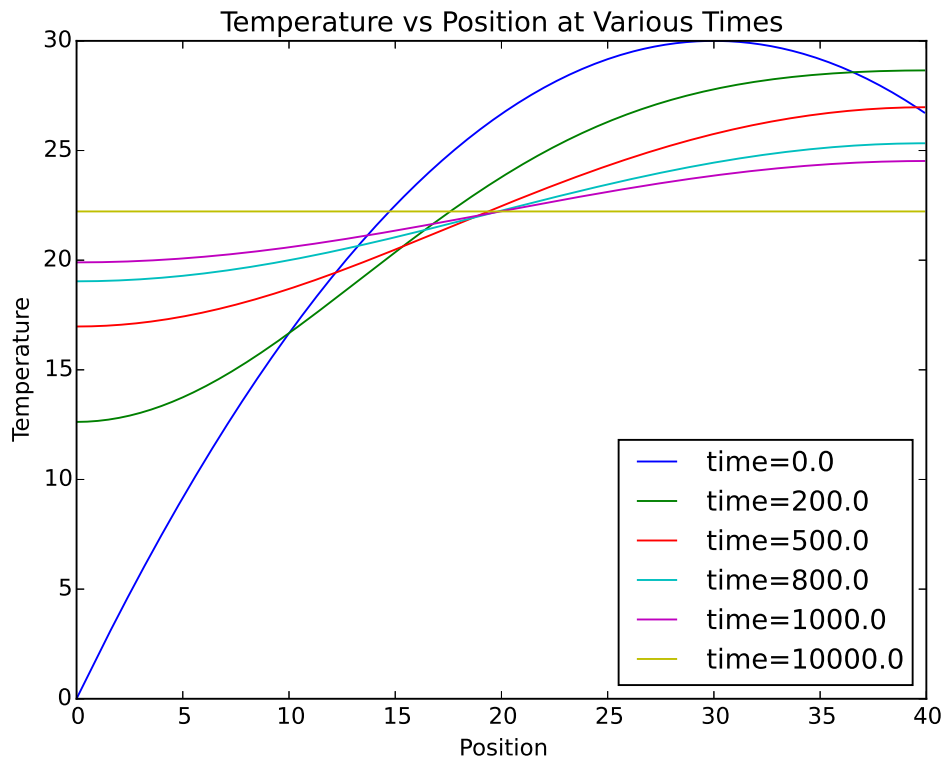
which works except for $n = 0$, for which we obtain $a_0 = 400/9$. Therefore we see

$$u(x, 0) = 200/9 + \sum_{n=1}^{\infty} \frac{160}{3} \frac{(-1)^{n+1} - 3}{n^2\pi^2} \cos(n\pi x/40).$$

We conclude that

$$u(x, t) = 200/9 + \sum_{n=1}^{\infty} \frac{160}{3} \frac{(-1)^{n+1} - 3}{n^2\pi^2} e^{-n^2\pi^2 t/6400} \cos(n\pi x/40).$$

- (b) Again, the exponential terms die off, so the steady state temperature is 200/9.
- (c) Plot at several times is included in the figure below.



Problem 3 *Nonhomogeneous Boundary Conditions*

Let an aluminum rod of length 20 cm be initially at the uniform temperature of 25 degrees C. Suppose that at time $t = 0$, the end $x = 0$ is cooled to 0 degrees C while the other end $x = 20$ is heated to 60 degrees C, and both are thereafter maintained at those temperatures.

- (a) Find the temperature $u(x, t)$. (Note: the initial condition $u(x, 0)$ does not satisfy the boundary conditions, which is fine since we are only asking the boundary conditions to be satisfied for $t > 0$)
- (b) What is the steady state temperature as $t \rightarrow \infty$?
- (c) Plot u vs. x for several values of t .
- (d) Determine how much time must elapse before the temperature at $x = 5$ comes within 1 degrees C of its steady state value.

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Solution 3.

- (a) We have Dirichlet boundary conditions, so this means that we should be thinking about a sine expansion. The first thing we should do is obtain the steady state solution $u_{\text{steady}}(x)$. Note that it is **not** time dependent (since it's steady state!!). Since it satisfies the heat equation, we know that $u''_{\text{steady}}(x) = 0$, and therefore $u_{\text{steady}}(x) = ax + b$ for some constants a and b . Now since $u_{\text{steady}}(0) = 0$ and $u_{\text{steady}}(20) = 60$, we can work out a and b , obtaining ($a = 3, b = 0$):

$$u_{\text{steady}}(x) = 3x.$$

Next, we should solve for the transient solution, which has homogeneous Dirichlet boundary conditions and initially satisfies $u_{\text{trans}}(x, 0) = 25 - u_{\text{steady}}(x) = 25 - 3x$. The sine expansion of this is given by

$$u_{\text{trans}}(x, 0) = \sum_{n=1}^{\infty} \frac{10}{n\pi} (5 + 7(-1)^n) \sin(n\pi x/20).$$

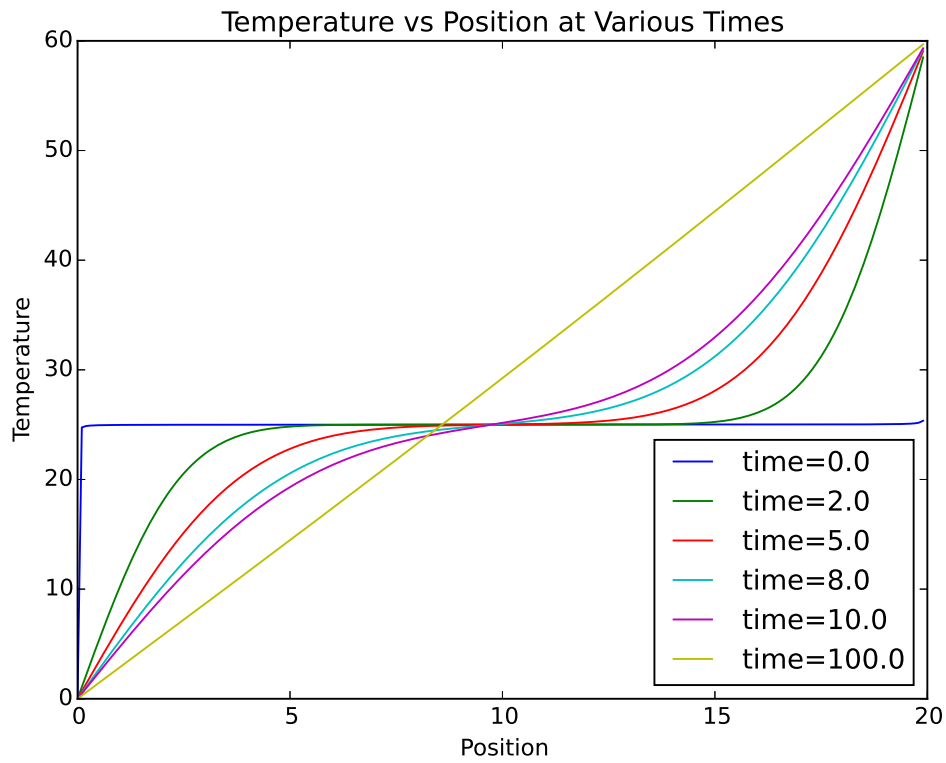
Then since $\alpha^2 = 0.86$ for aluminum, we find that the transient solution is:

$$u_{\text{trans}}(x, t) = \sum_{n=1}^{\infty} \frac{10}{n\pi} (5 + 7(-1)^n) e^{-n^2\pi^2(0.86)t/400} \sin(n\pi x/20).$$

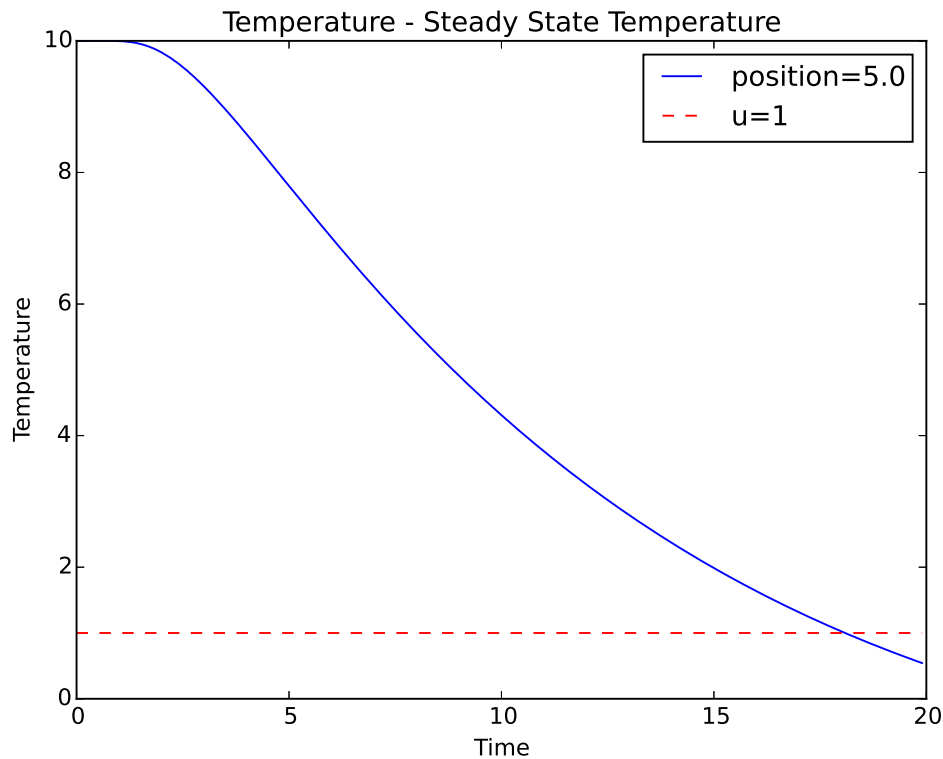
To get the solution to the inhomogeneous heat equation problem above, we need to add this transient solution to the steady state solution $u_{\text{steady}}(x)$. Therefore

$$u = u_{\text{trans}}(x, t) + u_{\text{steady}}(x, t) = 3x + \sum_{n=1}^{\infty} \frac{10}{n\pi} (5 + 7(-1)^n) e^{-n^2\pi^2(0.86)t/400} \sin(n\pi x/20).$$

- (b) As $t \rightarrow \infty$, the transient solution dies off, leaving behind only the steady state solution $u_{\text{steady}} = 3x$.
- (c) Plot at several times is included in the figure below.



(d) We are really asking how much time elapses until $|u_{\text{trans}}(5, t)| \leq 1$. We can determine this by plotting $u_{\text{trans}}(5, t)$ and determining what time it drops below 1 degree Celsius. A graph of the transient temperature is included below:



From the graph, we approximate $t = 18.18$ is about when $u(5, t)$ is within its steady state value.

Problem 4 *The Heat Equation in Two Dimensions*

We consider the two dimensional heat equation

$$u_t - \alpha^2(u_{xx} + u_{yy}) = 0.$$

- (a) Assume that u is of the form $u(x, y, t) = F(x)G(y)T(t)$, and show that the heat equation reduces to the system of three ordinary differential equations

$$\begin{cases} T'(t) + \lambda T = 0 \\ F''(x) + \frac{\lambda - \mu}{\alpha^2} F(x) = 0 \\ G''(y) + \frac{\mu}{\alpha^2} G(y) = 0 \end{cases}$$

for some constants λ and μ .

- (b) Assume that $u(x, y, t) = F(x)G(y)T(t)$ satisfies the heat equation above in the rectangular region $[0, L] \times [0, M]$ and also satisfies the Dirichlet boundary conditions

$$u(0, y, t) = 0, u(L, y, t) = 0, u(x, 0, t) = 0, u(x, M, t) = 0.$$

Find all possible functions $u(x, y, t)$ satisfying the above conditions. [Hint: they should be indexed by pairs of positive integers (m, n)]

- (c) Use (b) to find a solution to the two dimensional heat equation with Dirichlet boundary conditions

$$u_t - (u_{xx} + u_{yy}) = 0,$$

$$u(0, y, t) = 0, u(1, y, t) = 0, u(x, 0, t) = 0, u(x, 1, t) = 0,$$

with the initial condition that

$$u(x, y, 0) = \sin(3\pi x) \sin(2\pi y) + \sin(2\pi x) \sin(4\pi y).$$

Create a surface plots of your solution for several values of t .

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Solution 4.

- (a) Pluggin $u(x, y, t) = F(x)G(y)T(t)$ into the two dimensional heat equation, we obtain

$$F(x)G(y)T'(t) = \alpha^2(F''(x)G(y)T(t) + F(x)G''(y)T(t)).$$

Then dividing on both sides by $F(x)G(y)T(t)$, we obtain:

$$T'(t)/T(t) = \alpha^2(F''(x)/F(x) + G''(y)/G(y)).$$

The expression on the left hand side is a function of t only, while the expression on the right hand side is independent of t . Therefore both must be equal to a constant $-\lambda$:

$$T'(t)/T(t) = \alpha^2(F''(x)/F(x) + G''(y)/G(y)) = -\lambda.$$

Simplifying this, we get

$$T'(t)/T(t) = -\lambda$$

$$\alpha^2 F''(x)/F(x) + \alpha^2 G''(y)/G(y) = -\lambda.$$

Therefore

$$\alpha^2 F''(x)/F(x) = -\alpha^2 G''(y)/G(y) - \lambda.$$

Again the expression on the left hand side is a function of x only, and on the right we have a function of y only, and therefore both are equal to a constant μ :

$$\alpha^2 F''(x)/F(x) + \lambda = -\alpha^2 G''(y)/G(y) = \mu.$$

It follows that

$$\alpha^2 F''(x)/F(x) = \mu - \lambda$$

$$-\alpha^2 G''(y)/G(y) = -\mu.$$

Simplyfying things further we obtain the system of three ordinary differential equations listed in (a) above.

- (b) The Dirichlet boundary conditions result in boundary conditions on our various ODE's. In particular

$$F''(x) + \frac{\lambda - \mu}{\alpha^2} F(x) = 0, \quad F(0) = 0, \quad F(L) = 0$$

and also

$$G''(y) + \frac{\mu}{\alpha^2} G(y) = 0, \quad G(0) = 0, \quad G(M) = 0.$$

Then from our experience with boundary value problems, to get a nontrivial solution this says that $\mu = n^2\pi^2\alpha^2/M^2$ and that $\lambda - \mu = m^2\pi^2\alpha^2/L^2$ for some integers m and n , and in this case the solutions we obtain are

$$F(x) = A \sin(m\pi x/L), \quad G(y) = B \sin(n\pi y/M)$$

for some constants A and B . Then $\lambda = m^2\pi^2\alpha^2/L^2 + n^2\pi^2\alpha^2/M^2$, and therefore

$$T = C \exp\left(\frac{m^2\pi^2\alpha^2}{L^2}t + \frac{n^2\pi^2\alpha^2}{M^2}t\right).$$

Thus we obtain the solution

$$u(x, y, t) = F(x)G(y)T(t) = ABC \exp\left(\frac{m^2\pi^2\alpha^2}{L^2}t + \frac{n^2\pi^2\alpha^2}{M^2}t\right) \sin(m\pi x/L) \sin(n\pi y/M).$$

This motivates us to set

$$u_{mn}(x, y, t) = \exp\left(\frac{m^2\pi^2\alpha^2}{L^2}t + \frac{n^2\pi^2\alpha^2}{M^2}t\right) \sin(m\pi x/L) \sin(n\pi y/M).$$

Then by the superposition principle, any solution of the form

$$u(x, y, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} a_{mn} u_{mn}(x, y, t)$$

is a solution, for constants a_{mn} . In fact, all solutions may be written this way!

- (c) From part (b), we know to try to write

$$u(x, y, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} a_{mn} u_{mn}(x, y, t)$$

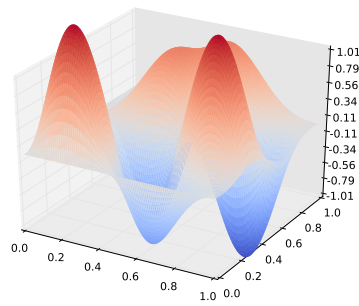
for some constants a_{mn} , where here $\alpha^2 = 1, L = 1, M = 1$. we must choose the constants so that

$$u(x, y, 0) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} a_{mn} u_{mn}(x, y, 0) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} a_{mn} \sin(m\pi x) \sin(n\pi y).$$

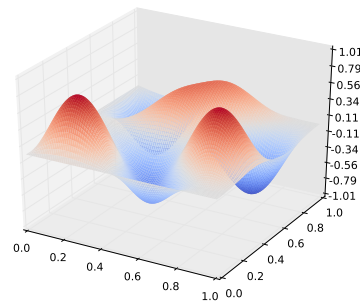
is equal to our initial condition. Looking at initial condition, this is easy! Just choose $a_{32} = 1, a_{24} = 1$ and $a_{mn} = 0$ otherwise. Thus:

$$u(x, y, t) = u_{32}(x, y, t) + u_{24}(x, y, t) = e^{-13\pi^2 t} \sin(3\pi x) \sin(2\pi y) + e^{-20\pi^2 t} \sin(2\pi x) \sin(4\pi y).$$

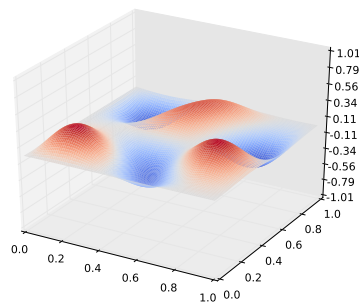
Plots at various times are included below:



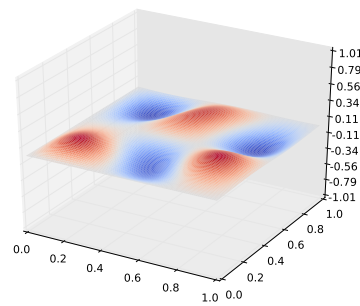
$t = 0.000$



$t = 0.005$



$t = 0.010$



$t = 0.015$

Problem 5 *The Heat Equation in Polar Coordinates*

We consider the two dimensional heat equation

$$u_t - \alpha^2(u_{xx} + u_{yy}) = 0.$$

(a) Show that using polar coordinates, (r, θ) , the heat equation becomes

$$u_t - \alpha^2 \left(u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} \right) = 0.$$

(b) Assume that u is of the form $u(r, \theta, t) = R(r)S(\theta)T(t)$, and show that the heat equation reduces to the system of three ordinary differential equations

$$\begin{cases} T'(t) + \lambda T = 0 \\ r^2 R''(r) + rR'(r) + \frac{1}{\alpha^2}(r^2\lambda - \mu)R = 0 \\ S''(\theta) + \frac{\mu}{\alpha^2}S(\theta) = 0 \end{cases}$$

for some constants λ and μ .

- (c) Explain why $\mu = n^2\alpha^2$ for some integer n . [Hint: remember that θ is the angle counter-clockwise from the x -axis].
- (d) Find the general solution to the above system of equations in the case that $\lambda = 0$ and $\mu = \alpha^2$. [Hint: to solve for $R(r)$, propose a solution of the form $R(r) = r^b$]

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Solution 5.

- (a) Note that $r^2 = x^2 + y^2$ and $\tan(\theta) = y/x$ and therefore

$$r_x = \cos(\theta), \quad r_y = \sin(\theta),$$

as well as

$$\theta_x = \frac{-\sin(\theta)}{r}, \quad \theta_y = \frac{\cos(\theta)}{r},$$

Then from the chain rule we have

$$\begin{aligned} \frac{\partial}{\partial x} &= r_x \frac{\partial}{\partial r} + \theta_x \frac{\partial}{\partial \theta} = \cos(\theta) \frac{\partial}{\partial r} - \frac{1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \\ \frac{\partial}{\partial y} &= r_y \frac{\partial}{\partial r} + \theta_y \frac{\partial}{\partial \theta} = \sin(\theta) \frac{\partial}{\partial r} + \frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \end{aligned}$$

Then we have that

$$\begin{aligned} \frac{\partial^2}{\partial x^2} &= \left(\frac{\partial}{\partial x} \right)^2 = \left[\cos(\theta) \frac{\partial}{\partial r} + \frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right]^2 \\ &= \left[\cos(\theta) \frac{\partial}{\partial r} \right]^2 + \left[\cos(\theta) \frac{\partial}{\partial r} \right] \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right] \\ &\quad + \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right] \left[\cos(\theta) \frac{\partial}{\partial r} \right] + \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right]^2 \end{aligned}$$

and also

$$\begin{aligned} \left[\cos(\theta) \frac{\partial}{\partial r} \right]^2 &= \cos^2(\theta) \frac{\partial^2}{\partial r^2} \\ \left[\cos(\theta) \frac{\partial}{\partial r} \right] \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right] &= \frac{-1}{r} \sin(\theta) \cos(\theta) \frac{\partial^2}{\partial r \partial \theta} + \frac{1}{r^2} \sin(\theta) \cos(\theta) \frac{\partial}{\partial \theta} \\ \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right] \left[\cos(\theta) \frac{\partial}{\partial r} \right] &= \frac{-1}{r} \sin(\theta) \cos(\theta) \frac{\partial^2}{\partial r \partial \theta} + \frac{1}{r} \sin^2(\theta) \frac{\partial}{\partial r} \\ \left[\frac{-1}{r} \sin(\theta) \frac{\partial}{\partial \theta} \right]^2 &= \frac{1}{r^2} \sin^2(\theta) \frac{\partial^2}{\partial \theta^2} + \frac{1}{r^2} \sin(\theta) \cos(\theta) \frac{\partial}{\partial \theta} \end{aligned}$$

Similarly,

$$\begin{aligned} \frac{\partial^2}{\partial y^2} &= \left(\frac{\partial}{\partial y} \right)^2 = \left[\sin(\theta) \frac{\partial}{\partial r} + \frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right]^2 \\ &= \left[\sin(\theta) \frac{\partial}{\partial r} \right]^2 + \left[\sin(\theta) \frac{\partial}{\partial r} \right] \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right] \\ &\quad + \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right] \left[\sin(\theta) \frac{\partial}{\partial r} \right] + \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right]^2 \end{aligned}$$

and also

$$\begin{aligned} \left[\sin(\theta) \frac{\partial}{\partial r} \right]^2 &= \sin^2(\theta) \frac{\partial^2}{\partial r^2}. \\ \left[\sin(\theta) \frac{\partial}{\partial r} \right] \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right] &= \frac{1}{r} \sin(\theta) \cos(\theta) \frac{\partial^2}{\partial r \partial \theta} - \frac{1}{r^2} \sin(\theta) \cos(\theta) \frac{\partial}{\partial \theta}. \\ \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right] \left[\sin(\theta) \frac{\partial}{\partial r} \right] &= \frac{1}{r} \sin(\theta) \cos(\theta) + \frac{1}{r} \cos^2(\theta) \frac{\partial}{\partial r}. \\ \left[\frac{1}{r} \cos(\theta) \frac{\partial}{\partial \theta} \right]^2 &= \frac{1}{r^2} \cos^2(\theta) \frac{\partial^2}{\partial \theta^2} - \frac{1}{r^2} \sin(\theta) \cos(\theta) \frac{\partial}{\partial \theta}. \end{aligned}$$

Adding all of this together, we obtain:

$$\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}.$$

Applying this to the function u , we get

$$u_{xx} + u_{yy} = u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta}.$$

Substituting this into the heat equation leads to

$$u_t - \alpha^2 \left(u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} \right) = 0.$$

(b) Putting $u(r, \theta, t) = R(r)S(\theta)T(t)$ into the above expression, we obtain:

$$R(r)S(\theta)T'(t) = \alpha^2 \left(R''(r)S(\theta)T(t) + \frac{1}{r} R'(r)S(\theta)T(t) + \frac{1}{r^2} R(r)S''(\theta)T(t) \right).$$

Dividing everything by $R(r)S(\theta)T(t)$, we obtain

$$T'(t)/T(t) = \alpha^2 \left(R''(r)/R(r) + \frac{1}{r} R'(r)/R(r) + \frac{1}{r^2} S''(\theta)/S(\theta) \right).$$

The expression on the left hand side is a function of t only, while the right hand side is a function of r and θ only, and therefore both are equal to an arbitrary constant $-\lambda$. This leads to

$$T'(t)/T(t) = -\lambda$$

and also

$$\alpha^2(R''(r)/R(r) + \frac{1}{r}R'(r)/R(r) + \frac{1}{r^2}S''(\theta)/S(\theta)) = -\lambda.$$

Then after an algebraic manipulation,

$$\alpha^2r^2R''(r)/R(r) - \alpha^2rR'(r)/R(r) - r^2\lambda = S''(\theta)/S(\theta).$$

Therefore again we have that both are equal to a constant $-\mu$:

$$\alpha^2r^2R''(r)/R(r) - \alpha^2rR'(r)/R(r) - r^2\lambda = -\mu$$

and also

$$S''(\theta)/S(\theta) = -\mu.$$

After some simplification, this reduces to the expression in (b) above.

- (c) Since (r, θ) and $(r, \theta + 2\pi)$ both refer to the same coordinate in polar coordinates, the function S should satisfy $S(\theta) = S(\theta + 2\pi)$. Moreover, we know that S satisfies $S''(\theta) + (\mu/\alpha^2)S(\theta) = 0$. Depending on the value of μ/α^2 , the solutions to this are either exponentials, polynomials or trig. functions. Since we require S to be periodic, we want them to be trig functions, and therefore we need μ/α^2 to be positive. In this case, the general solution to the differential equation is

$$S(\theta) = A \cos(\sqrt{\mu}\theta/\alpha) + B \sin(\sqrt{\mu}\theta/\alpha).$$

Now we need S to be 2π periodic, and therefore we need $\sqrt{\mu}/\alpha = n$ for some integer n . Therefore $\mu = n^2\alpha^2$.

- (d) Since $\lambda = 0$, the solution for T is $T = E$ for some constant E . Moreover, the general solution for S is

$$S(\theta) = A \cos(\theta) + B \sin(\theta)$$

Finally, we propose a solution $R(r) = r^b$ to the equation $r^2R''(r) + rR'(r) - R(r) = 0$, and therefore $b(b-1)r^{b-2}r^b + br^b - r^b = 0$. This leads to $b^2 - 1 = 0$, and therefore $b = \pm 1$. Thus we obtain two solutions: r and r^{-1} . The general solution is therefore

$$R(r) = Cr + Dr^{-1}$$

. Putting this all together, we obtain

$$u(r, \theta, t) = R(r)S(\theta)T(t) = (Cr + Dr^{-1})(A \cos(\theta) + B \sin(\theta))E.$$

Since A, B, C, D, E are all arbitrary constants, we can rewrite this as:

$$u(r, \theta, t) = C_1r \cos(\theta) + C_2r \sin(\theta) + C_3r^{-1} \cos(\theta) + C_4r^{-1} \sin(\theta).$$